CÔTE D'AZUR

## Quantum turbulence: From the Kolmogorov cascade to sound emission, passing by Kelvin waves and vortex reconnections.

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Le calcul intensif au service de la connaissance


## Superfluids

## Superfluid ${ }^{4} \mathrm{He}$

$$
T \sim 2 K
$$

## They have no viscosity !

Helium phase diagram


Picture from: Low Temperature Laboratory, Aalto University webpage

## Landau-Tiszla description of superfluid helium

## Two immiscible fluids: <br> $\rightarrow$ normal (viscous) fluid of density $\rho_{\mathrm{n}}$ <br> $\rightarrow$ superfluid of density $\rho_{\mathrm{s}}$ <br> $$
\begin{aligned} & \rho=\rho_{\mathrm{n}}+\rho_{\mathrm{s}} \\ & \mathbf{P}=\rho_{\mathrm{n}} \mathbf{v}_{\mathrm{n}}+\rho_{\mathrm{s}} \mathbf{v}_{\mathrm{s}} \end{aligned}
$$



Today's talk

## Length scales of superfluid turbulence

energy
injection
$\sim m$

Classical (Kolmogorov) turbulence


SHREK (France)
inter-vortex distance
$\sim 10^{-5} \mathrm{~m}$

Kelvin wave cascade \& vortex reconnections
coherence length vortex core size
$\xi$


Experiments: Maurer et al. (1998), Salort et al. (2010), Tang et al. (2021),
Simulations in GP: Nore et al. (1997), Kobayashi et al. (2005),
Simulations in vortex-filament method: Baggaley et al. (2012),

G. Bewley et al. Nature 2006.


## Quantum vortices and turbulence

At "zero-temperature", a superfluid has no viscosity
Compressible fluid (and dispersive)
Described by a complex order parameter (wave function)
Quantum vortices (filaments) are naturally present in turbulent states


## Quantum vortices


-Finite core-size
Continuous circulation


- Topological defects

Quantised circulation

# Modeling superfluid helium <br> Multi-scale physics 

## Scales

vortex core size

mean inter-vortex distance
$\ell$

## Modeling superfluid helium

Multi-scale physics

## Scales



## Classical fluid (Navier-Stokes)

# Modeling superfluid helium <br> Multi-scale physics 

## Scales

vortex core size
$T=0$
mean inter-vortex distance
$\ell$

Gross-Pitaevskii based model

Classical fluid (Navier-Stokes)

## The Gross-Pitaevskii equation <br> Modelling low-temperature superfluids

$$
i \hbar \frac{\partial \psi}{\partial t}=-\frac{\hbar^{2}}{2 m} \nabla^{2} \psi+g|\psi|^{2} \psi, \quad g=\frac{4 \pi a \hbar^{2}}{m}
$$

Linearising about a flat state:

$$
\psi=A_{0} e^{-i \frac{\mu}{\hbar} t}+\delta \psi
$$

Bogoliubov dispersion relation:

$$
\omega(k)=c k \sqrt{1+\frac{1}{2} \xi^{2} k^{2}}
$$

$$
\omega(k)=\sqrt{\frac{g\left|A_{0}\right|^{2}}{m} k^{2}+\frac{\hbar^{2}}{4 m^{2}} k^{4}} .
$$

$$
\begin{array}{cl}
\text { Speed of sound } & c=\sqrt{g\left|A_{0}\right|^{2} / m} \\
\text { Coherence length } & \xi=\sqrt{\hbar^{2} / 2 m\left|A_{0}\right|^{2} g}
\end{array}
$$

## Hydrodynamics?

$i \hbar \frac{\partial \psi}{\partial t}=-\frac{\hbar^{2}}{2 m} \nabla^{2} \psi+g|\psi|^{2} \psi$,

$$
\begin{aligned}
\text { Speed of sound } & c=\sqrt{g\left|A_{0}\right|^{2} / m} \\
\text { Coherence length } & \xi=\sqrt{\hbar^{2} / 2 m\left|A_{0}\right|^{2} g}
\end{aligned}
$$

Madelung transformation

$$
\psi(\mathbf{x}, t)=\sqrt{\frac{\rho(\mathbf{x}, t)}{m}} \exp \left[i \frac{m}{\hbar} \phi(\mathbf{x}, t)\right]=\sqrt{\frac{\rho(\mathbf{x}, t)}{m}} \exp \left[i \frac{\phi(\mathbf{x}, t)}{\sqrt{2} c \xi}\right]
$$

density of particles

$$
\frac{\partial \rho}{\partial t}+\boldsymbol{\nabla} \cdot(\rho \nabla \phi)=0
$$

$$
\frac{\partial \phi}{\partial t}+\frac{1}{2}(\nabla \phi)^{2}=c^{2}(1-\rho)+c^{2} \xi^{2} \frac{\Delta \sqrt{\rho}}{\sqrt{\rho}} .
$$

$$
\rho=|\psi|^{2} \quad \mathbf{v}=\nabla \phi \quad \text { is a potential flow }
$$

## Quantum vortices

$\mathbf{v}=\nabla \phi$ is a potential flow but:
Vortices are topological defects: $\psi(\mathbf{x})=\mathbf{0}$

$$
\Gamma=\oint_{\mathcal{C}} \nabla \phi \cdot \mathrm{d} \ell=\phi^{+}-\phi^{-}
$$

$$
\Gamma=n \frac{h}{m}=n 2 \pi \sqrt{2} c \xi, \quad \text { with } n \in \mathbb{Z}
$$

Points in $2 D$ and lines in $3 D$


$$
\mathbf{v} \sim \frac{1}{r} \Rightarrow \nabla \times \mathbf{v} \sim \delta(\mathbf{r})
$$

$$
\mathbf{w}(\mathbf{x})=\nabla \times \mathbf{v}=\frac{h}{m} \oint \delta(\mathbf{x}-\mathbf{s}(\ell)) \frac{\mathrm{d} \mathbf{s}(\ell)}{\mathrm{d} \ell} \mathrm{~d} \ell
$$



## Quantum vortices

## $\mathbf{v}=\nabla \phi$ is a potential flow but:

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$$
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$$

## Gross-Pitaevskii model

Collection of "ideal" vortex filaments:

- Velocity field $\vec{v}=\frac{\Gamma}{2 \pi r} \hat{\theta}$
- Core size $\xi$ (small)


Density $\rho$
Euler equation + small-scale dispersion regularisations

$$
\mathbf{w}(\mathbf{x})=\nabla \times \mathbf{v}=\frac{h}{m} \oint \delta(\mathbf{x}-\mathbf{s}(\ell)) \frac{\mathrm{d} \mathbf{s}(\ell)}{\mathrm{d} \ell} \mathrm{~d} \ell
$$



## Vortex fllament method

$$
\text { At } T=0
$$

$$
\mathbf{v} \sim \frac{1}{r} \Rightarrow \nabla \times \mathbf{v} \sim \delta(\mathbf{r})
$$

## Collection of vortex

 filaments $\mathscr{C}$$$
\text { Vorticity field } \omega_{\mathbf{s}}(\mathbf{x})=\Gamma \oint_{\mathscr{C}} \delta(\mathbf{x}-\mathbf{s}(\zeta)) \frac{d \mathbf{R}}{d \zeta} d \zeta
$$

Velocity field

$$
\mathbf{v}_{\mathbf{s}}(\mathbf{x})=-\frac{1}{4 \pi} \int \frac{(\mathbf{x}-\mathbf{y})}{|\mathbf{x}-\mathbf{y}|^{3}} \times \omega(\mathbf{y}) d^{3} \mathbf{y}=-\frac{\Gamma}{4 \pi} \oint_{\mathscr{C}} \frac{\mathbf{x}-\mathbf{s}(\zeta)}{|\mathbf{x}-\mathbf{s}(\zeta)|^{3}} \times \frac{d \mathbf{R}}{d \zeta} d \zeta
$$

$$
\frac{d \mathbf{R}(\zeta, \mathbf{t})}{d t}=\mathbf{v}_{\mathbf{s}}(\mathbf{R}(\zeta, t))
$$

## Length scales of

## Today's talk

energy
injection
$\sim m$

Classical (Kolmogorov) turbulence
.......

inter-vortex distance
$\sim 10^{-5} m$

Kelvin wave cascade
\& vortex reconnections
coherence length
vortex core size


## Kelvin waves

## XXIV. Vibrations of a Columnar Vortex. By Sir William Thomson*.

$T$ HIIS is a case of fluid-motion, in which the stream-lines are approximately circles, with their centres in one line (the axis of the vortex) and the velocities approximately constant, and approximately equal at equal distances from the axis. As a preliminary to treating it, it is convenient to ex-


Sir William Thomson (I880) XXIV. Philosophical Magazine Series 5, 10:6I, I55-I68,
Take the incompressible Euler's equations

$$
\begin{gathered}
\mathbf{v}_{0}(r, \theta, z)=\frac{\alpha(r)}{r} \hat{\theta} \quad \text { and } \quad p(r, \theta, z)=p_{0}(r)=\rho_{0} \int_{a_{0}}^{r} \frac{\alpha(s)^{2}}{s^{3}} \mathrm{~d} s \\
\qquad=\oint_{\mathcal{C}} \mathbf{v} \cdot \mathrm{d} \ell=2 \pi \alpha(r) \\
\text { Kelvin Waves: } \quad \mathbf{v}=\mathbf{v}_{0}+\delta \mathbf{v}+\ldots
\end{gathered}
$$

## Kelvin waves

$\delta v(\theta, z) \sim \cos (k z) \sin (n \theta-\omega t)$

$$
\omega_{n}^{ \pm}(k)=\frac{\Gamma}{2 \pi a_{0}^{2}}\left(n \pm \sqrt{n+\frac{a_{0}|k| K_{n-1}\left(a_{0}|k|\right)}{K_{n}\left(a_{0}|k|\right)}}\right)
$$



$n=1$ and $k a_{0} \ll 1$

$$
\left.\omega^{-}(k)=-\frac{\Gamma}{8 \pi} k^{2} \log \frac{1}{a_{0}|k|}+b\right), \text { with } b=\log 2-\gamma_{\mathrm{E}}
$$

## Vortex excitations in superfluids (GP)

Hydrodynamics
P.H. Roberts. Proc. Royal Society of London A:(2003)

$$
k \ll \frac{1}{\xi}
$$

$$
\begin{aligned}
\Omega_{\mathrm{KW}}\left(k a_{0} \rightarrow 0\right) & =-\frac{\Gamma}{4 \pi} k^{2}\left(\ln \frac{2}{a_{0}|k|}-\gamma_{\mathrm{E}}\right) \\
a_{0} & =1.265 \xi
\end{aligned}
$$

GP numerical simulations

U. Giuriato, G. Krstulovic and S. Nazarenko. Phys. Rev. Research (2020)

## Vortex excitations in superfluids (LIA)

## Local Induced Approximation (LIA)



$$
\dot{\mathbf{s}}=\frac{\Gamma \Lambda}{4 \pi R} \widehat{\mathbf{b}}=\frac{\Gamma \Lambda}{4 \pi} \mathbf{s}^{\prime} \times \mathbf{s}^{\prime \prime}
$$



## Small amplitudes Kelvin waves

$$
\begin{aligned}
& s(z, t)=X(z, t)+i Y(z, t) \\
& i \Gamma \dot{s}=\frac{\delta H_{\mathrm{LIA}}}{\delta s^{*}}=-\frac{\Gamma^{2} \Lambda}{4 \pi} \frac{\partial^{2} s}{\partial z^{2}}, \quad \text { with } \quad H_{\mathrm{LIA}}=\frac{\Gamma^{2} \Lambda}{4 \pi} \int\left|\frac{\partial s}{\partial z}\right|^{2} \mathrm{~d} z \\
& \omega_{\mathrm{LIA}}(k)=-\frac{\Gamma \Lambda}{4 \pi} k^{2}
\end{aligned}
$$

## Kelvin-wave cascade

## Vortex filament model

Biot-Savart description of a perturbed straight vortex
[Sonin 87 - Svistunov 95]


$$
\begin{aligned}
& s(z, t)=X(z, t)+i Y(z, t) \\
& i \Gamma \dot{s}(z)=\frac{\delta H_{\mathrm{NL}}}{\delta s^{*}(z)}, \quad H_{\mathrm{NL}}=\frac{\Gamma^{2}}{4 \pi} \int \frac{1+\mathcal{R} e\left[s^{\prime *}\left(z_{1}\right) s^{\prime}\left(z_{2}\right)\right]}{\sqrt{\left(z_{1}-z_{2}\right)^{2}+\left|s\left(z_{1}\right)-s\left(z_{2}\right)\right|^{2}}} \mathrm{~d} z_{1} \mathrm{~d} z_{2}
\end{aligned}
$$

Small amplitude waves:

$$
\begin{gathered}
H_{\mathrm{NL}}=\sum_{k} \omega_{k}\left|s_{k}\right|^{2}+H_{4}+H_{6}+\ldots \\
\quad \omega_{k}=-\frac{\Gamma}{4 \pi} k^{2}(\log (k \ell)-\Lambda)
\end{gathered}
$$

## Vortex motion and Kelvin wave cascade



## Wave turbulence predictions



Kozik-Svistunov (2004) :
(6 waves)

$$
\begin{gathered}
E_{\mathrm{KS}}(k)=C_{\mathrm{KS}} \frac{\Lambda \kappa^{7 / 5} \epsilon^{1 / 5}}{k^{7 / 5}} \\
\Lambda=\ln (\ell / a)
\end{gathered}
$$

a (wave-turbulence) controversy!!

L'vov-Nazarenko (2010):
(effective 4 wave theory )

$$
\begin{gathered}
E_{\mathrm{LN}}(k)=C_{\mathrm{LN}} \frac{\Lambda \kappa \epsilon^{1 / 3}}{\Psi^{2 / 3} k^{5 / 3}}, \quad \Psi \equiv \frac{8 \pi E}{\Lambda \kappa^{2}} \\
C_{\mathrm{LN}}=0.304
\end{gathered}
$$

- Kivotides, Vassilicos, Samuel, Barenghi PRL 200I
-E. Kozik \& B. Svistunov. PRL 2004
-L'vov \& Nazarenko JETP 2010
-Boué et al PRB 20II
-Laurie and Baggaley PRE 2014
-many others works....


## Kelvin-wave cascade

## Numerical simulations

## We consider a perturbed straight vortex:



Biot-Savart dynamics:

$$
\begin{aligned}
i \Gamma \dot{s}(z) & =\frac{\delta H_{\mathrm{NL}}}{\delta s^{*}(z)}, \quad H_{\mathrm{NL}}=\frac{\Gamma^{2}}{4 \pi} \int \frac{1+\mathcal{R} e\left[s^{\prime *}\left(z_{1}\right) s^{\prime}\left(z_{2}\right)\right]}{\sqrt{\left(z_{1}-z_{2}\right)^{2}+\left|s\left(z_{1}\right)-s\left(z_{2}\right)\right|^{2}}} \mathrm{~d} z_{1} \mathrm{~d} z_{2} \\
\dot{\mathbf{s}}(\zeta) & =\frac{\Gamma}{4 \pi} \oint \frac{\mathrm{~d} \mathbf{s}\left(\zeta^{\prime}\right) \times\left(\mathbf{s}(\zeta)-\mathbf{s}\left(\zeta^{\prime}\right)\right)}{\mid \mathbf{s}(\zeta)-s\left(\left.\zeta^{\prime}\right|^{3}\right.}
\end{aligned}
$$

Non-local equation, needs to be regularised, dissipation is added in an ad-hoc manner

Gross-Pitaevskii dynamics: $\quad i \hbar \frac{\partial \psi}{\partial t}=-\frac{\hbar^{2}}{2 m} \nabla^{2} \psi+g|\psi|^{2} \psi$,
3D PDE but everything is regular. Effective dissipation is provided by acoustic emission.
One gets $\psi(x, y, z, t)$ but we need a filament $\mathbf{s}(z, t)$ !

## Tracking vortices

G. Krstulovic

PRE 86, $055301(\mathrm{R}),(20 \mathrm{I} 2)$


Isosurface



Tracked lines

- Highly accurate (spectral precision)
- Geometry independent
- Arbitrary number of objects
A.Villois, G. Krstulovic, D. Proment and H. Salman. J. Phys. A (2016)


## Superfluid turbulence

energy
injection

Classical (Kolmogorov)
turbulence

Kelvin wave cascade vortex core
inter-vortex distance
size sound emission
$\rightarrow \ell_{I}$
size

## Is the Kelvin wave cascade relevant for a

## turbulent tangle?



## Quantum turbulence

Kelvin waves in a turbulent tangle


## Quantum turbulence

## Kelvin wave cascade



## Quantum turbulence

## Strong turbulence

Kolmogorov scaling for the energy spectrum (K41)

$$
\begin{gathered}
E(k)=C_{\mathrm{K}} \epsilon^{2 / 3} k^{-5 / 3} \\
k_{0} \ll k \ll k_{\ell}
\end{gathered}
$$

## Weak wave turbulence

Kelvin wave scaling for the energy spectrum

$$
\begin{gathered}
E(k) \sim \kappa \epsilon^{1 / 3} \ell^{-4 / 3} k^{-5 / 3} \\
k_{\ell} \ll k \ll k_{\xi}
\end{gathered}
$$

Non-local high-order nonlinearity GP


Simultaneous observation of two cascades

## Quantum turbulence

|  | Initial condition |  | Turbulence |  |
| :---: | :---: | :---: | :---: | :---: |
|  | $k_{0}$ | $L / \xi$ | $\epsilon$ | $\ell / L$ |
| $-=-=-$ | 2 | 341 | 0.01 | 0.412 |
| - | 2 | 171 | 0.01 | 0.494 |
| - | 2 | 341 | 0.01 | 0.255 |
| - | 3 | 341 | 0.02 | 0.235 |
| - | 4 | 341 | 0.03 | 0.227 |
| - | 2 | 683 | 0.01 | 0.139 |$\} \quad$ Local

Energy spectrum compensated by Kolmogorov spectrum


Energy spectrum compensated by Kelvin wave spectrum


## Vortex reconnections

Experiments in water


Kleckner \& Irvine. Nature
Phys. 2013

Numerical simulations of classical fluids


Navier-Stokes equations

Numerical simulations of superfluids


Gross-Pitaevskii model

Ideal for a theoretical description! $\quad \xi \ll R$

## Minimal vortex distance

Dimensional analisys: $\quad \xi$ :vortex core size $\quad \delta$ :minimal distance $\quad R$ :system size

$\xi \ll \delta \ll R$ between vortices

$$
\delta^{ \pm}(t)=A^{ \pm}\left|\Gamma\left(t-t_{r}\right)\right|^{1 / 2}
$$

## Minimal vortex distance



## Minimal vortex distance

| Analytic calculations | $\xi$ : vortex core size | $\delta$ : minimal distance | $R$ : system size |
| :---: | :---: | :---: | :---: |
| $\Gamma=\oint_{L^{2}} \mathbf{v} \cdot \mathrm{~d} \ell$ |  |  |  |
| $[\Gamma]=\frac{L^{-}}{T}$ | $\xi \ll \delta \ll L$ |  |  |
|  |  | $\delta^{ \pm}(t)=A^{ \pm}\left\|\Gamma\left(t-t_{r}\right)\right\|^{1 / 2}$ |  |



Previous works reported different exponents:

Zuccher et al Phys Fluids (20|3)
Allen et al. PRA (20|4)
Rorai et al.JFM (20|6)

## Quantum vortex filaments

4 study cases


Tangle
d.1)
d.2)


## Separation rates

a.1)
b.1)
c.1)
d.1)
,

a.2)

b.1)

c.1) Trefoil knot
d.1)

b.2)




## Geometry of reconnections



$$
\begin{aligned}
& \text { 米 } A^{+} \geq A^{-} \text {: vortices separate } \\
& \text { faster than they approach }
\end{aligned}
$$

Linear (Schrodinger equation) theory:
粦 $A^{+} / A^{-}$controls curvature, torsion and approach angle:

$$
\frac{A^{+}}{A^{-}}=\cot \left(\frac{\phi^{+}}{2}\right)
$$



$$
\frac{A^{+}}{A^{-}} \geqq 1 \Longleftrightarrow \phi^{+} \Longleftarrow \frac{\pi}{2}
$$

$$
\delta^{ \pm}(t)=A^{ \pm}\left|\Gamma\left(t-t_{r}\right)\right|^{1 / 2}
$$

## Quantum vortex reconnections

Physical
REVIEW
LETTERS





More than 40 reconnections


# QUANTUM VORTEX RECONNECTIONS Sound emission and irreversibility 

## Quantum vortex reconnections



Vortex separate faster than they approach

Matching theory


Analytical theory based on conservation of momentum and energy

## Quantum vortex reconnections

A more general ansatz

$$
\psi_{r}(x, y, z)=z+\frac{\gamma}{a}\left(x^{2}+y^{2}\right)+i\left(a z+\beta x^{2}-y^{2}\right)
$$


torsion

$\psi(\mathbf{x}, t)=e^{i \frac{1}{2}\left(t-t_{r}\right) \nabla^{2}} \psi_{r}(\mathbf{x})$

$$
\psi(\mathbf{x}, t)=0
$$

$$
\mathbf{R}_{1,2}^{-}(s, t)=\left(s, \pm \sqrt{\frac{\left(t_{r}-t\right)\left(a^{2}(1-\beta)-2 \gamma\right)+a s^{2}(\beta-\gamma)}{a(\gamma+1)}},\right.
$$

$$
\left.\frac{\left(t-t_{r}\right)\left(a^{2}(\beta-1)-2 \gamma^{2}\right)-a \gamma(\beta+1) s^{2}}{(\gamma+1) a^{2}}\right)
$$

$$
\delta^{ \pm}(t)=\left|\mathbf{R}_{1}^{ \pm}(0, t)-\mathbf{R}_{2}^{ \pm}(0, t)\right|=\sqrt{2 \pi} A^{ \pm}\left|t-t_{r}\right|^{1 / 2}
$$

time

$$
\begin{gathered}
\mathbf{P}^{-}=\frac{\Gamma}{2} \oint \mathbf{R}^{-} \times \mathrm{d} \mathbf{R}^{-} \\
E_{\text {kin }}^{-} \propto \Gamma^{2} \oint\left|\mathrm{~d} \mathbf{R}^{-}\right|
\end{gathered}
$$



Momentum

$$
\mathbf{P}^{+}=\frac{\Gamma}{2} \oint \mathbf{R}^{+} \times \mathrm{d} \mathbf{R}^{+}
$$

Energy

$$
E_{\text {kin }}^{+} \propto \Gamma^{2} \oint\left|\mathrm{~d} \mathbf{R}^{+}\right|
$$

$$
\Delta \mathbf{P}=\mathbf{P}^{+}-\mathbf{P}^{-}
$$

$$
\Delta E_{\mathrm{kin}}=E_{\mathrm{kin}}^{+}-E_{\text {kin }}^{-}
$$

$\frac{A^{+}}{A^{-}}=\sqrt{\frac{1+\gamma}{\beta-\gamma}}$

## Quantum vortex reconnections

## Numerical measurements



## QuantumVIW project



## Kelvin waves

## Vortex reconnections



